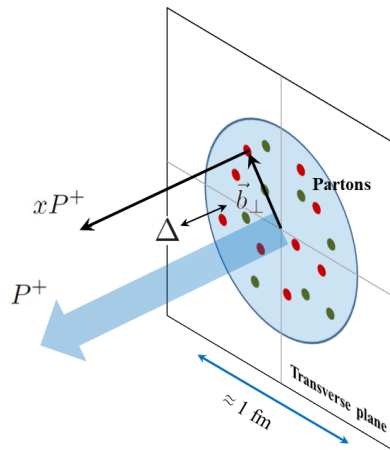




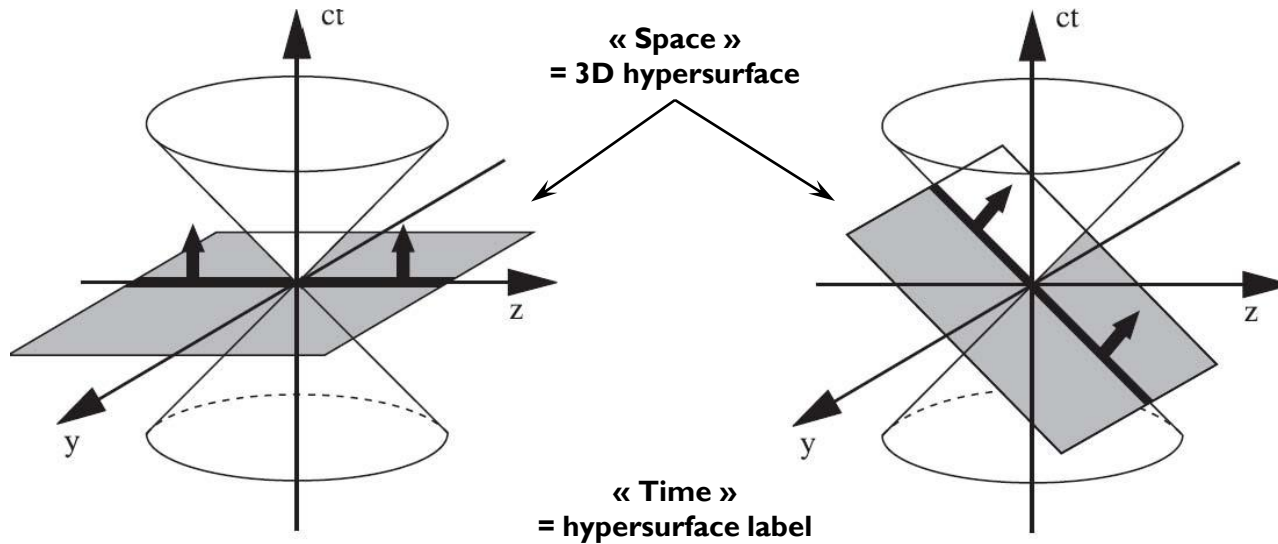
Hadron structure with GPDs (2/5)

Cédric Lorcé

Light-front picture

Forms of dynamics

Space-time foliation



Light-front components

$$a^{\pm} = \frac{1}{\sqrt{2}}(a^0 \pm a^3)$$

Instant form dynamics

Time	x^0
Space	$\vec{x}$
Energy	p^0
Momentum	$\vec{p}$

Light-front form dynamics

$$x^+$$

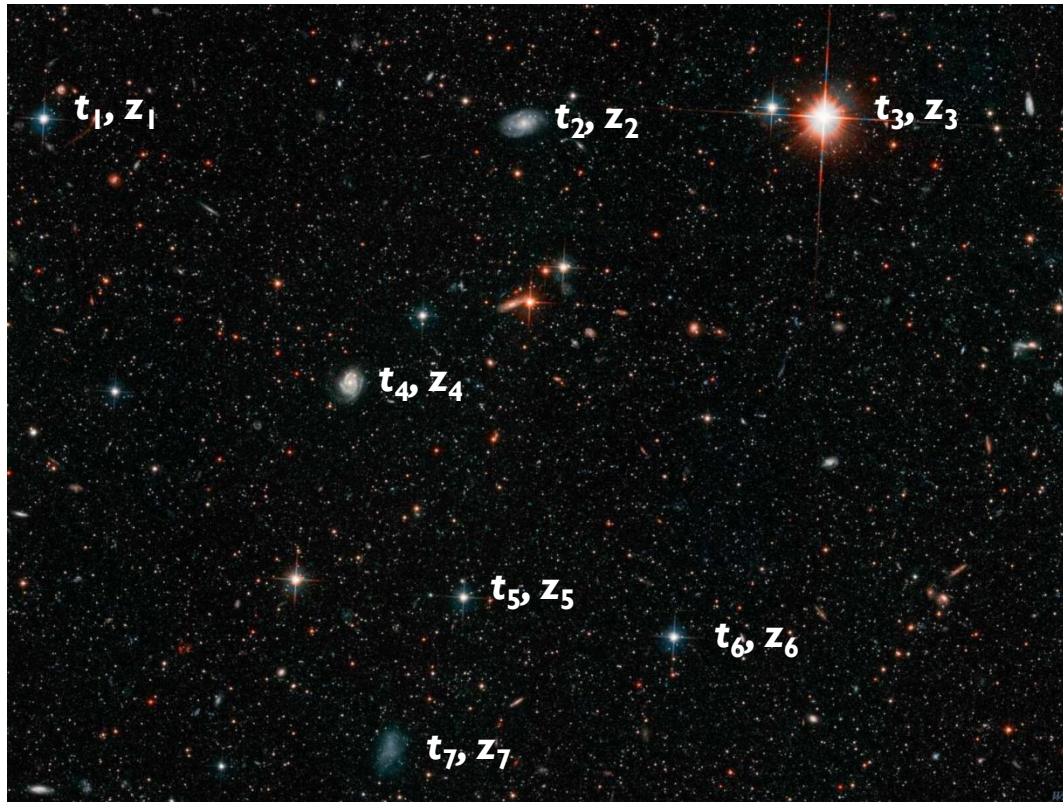
$$\vec{x}_{\perp}, x^{-}$$

$$p^{-}$$

$$\vec{p}_{\perp}, p^{+}$$

Instant form vs light-front form

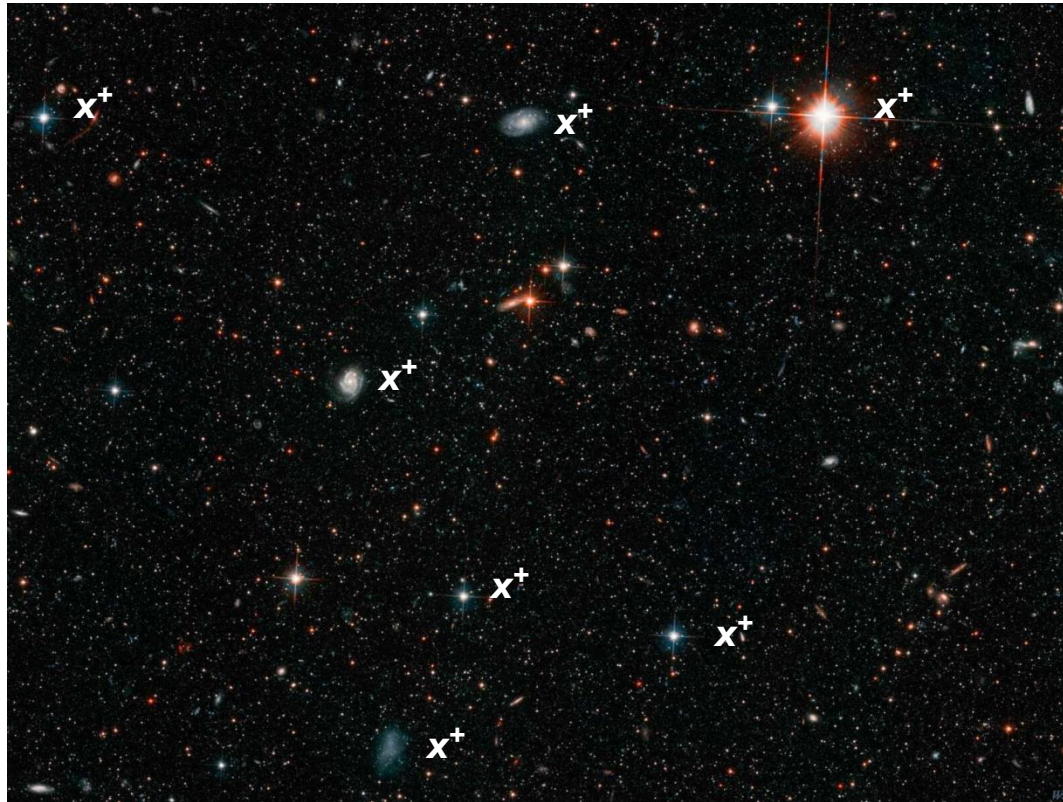
Ordinary point of view



The farther away, the earlier

Instant form vs light-front form

Light front point of view $x^+ = \frac{1}{\sqrt{2}}(t + z)$

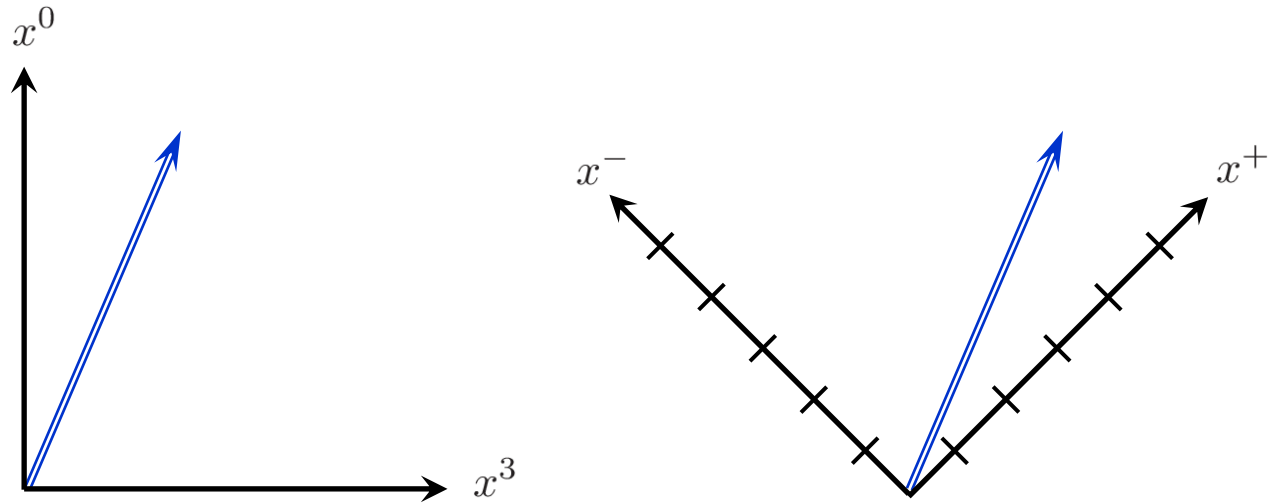


Change in synchronization protocol !

[Freese, Miller, PRD107 (2023) 7, 074036]
[Freese, Miller, PRD108 (2023) 9, 094026]

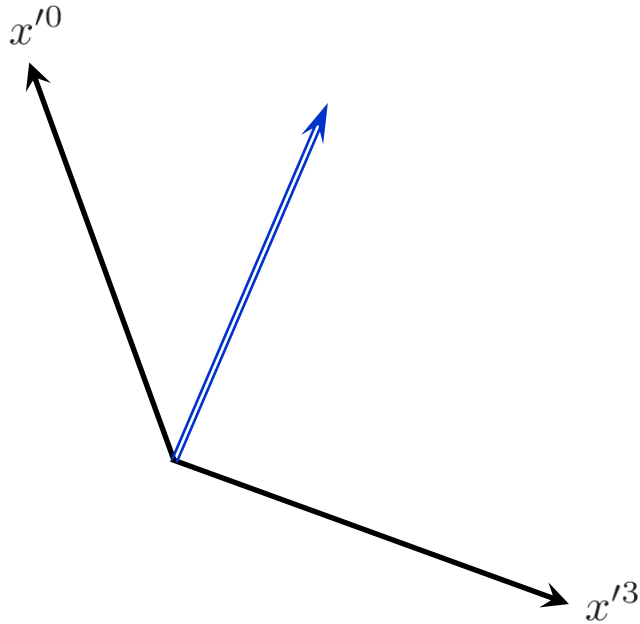
Instant form vs light-front form

Initial frame

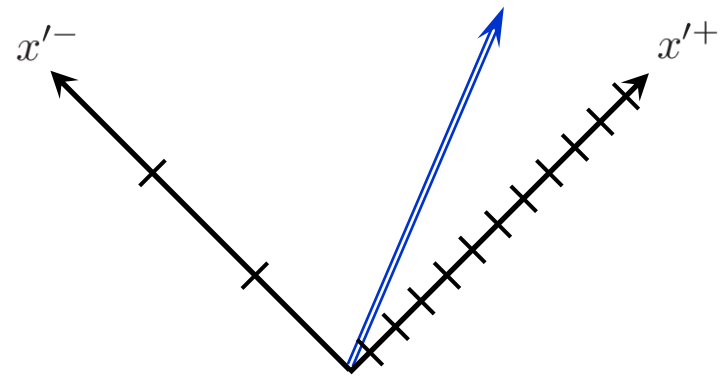


Instant form vs light-front form

Boosted frame



$$x'^0 = \frac{x^0 + \beta x^3}{\sqrt{1 - \beta^2}}$$
$$x'^3 = \frac{x^3 + \beta x^0}{\sqrt{1 - \beta^2}}$$

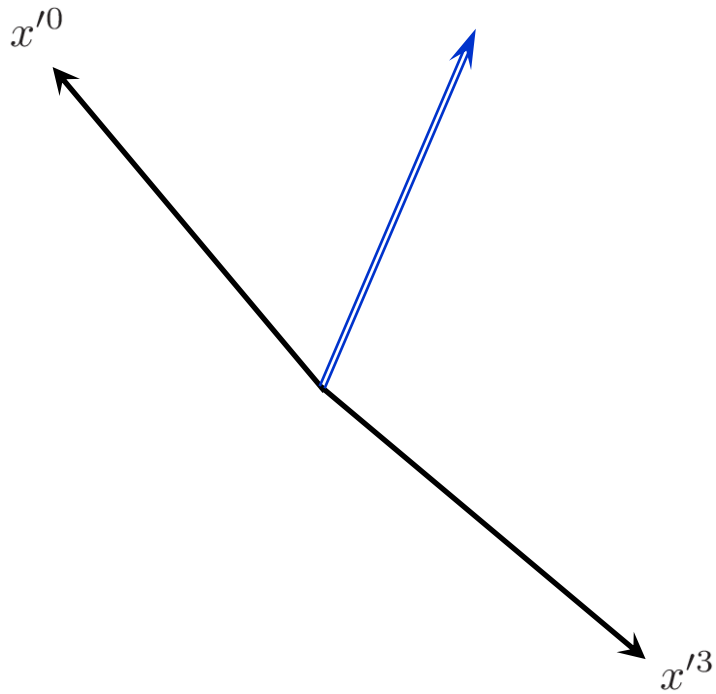


$$x'^+ = \sqrt{\frac{1 + \beta}{1 - \beta}} x^+ = e^{\eta} x^+$$
$$x'^- = \sqrt{\frac{1 - \beta}{1 + \beta}} x^- = e^{-\eta} x^-$$

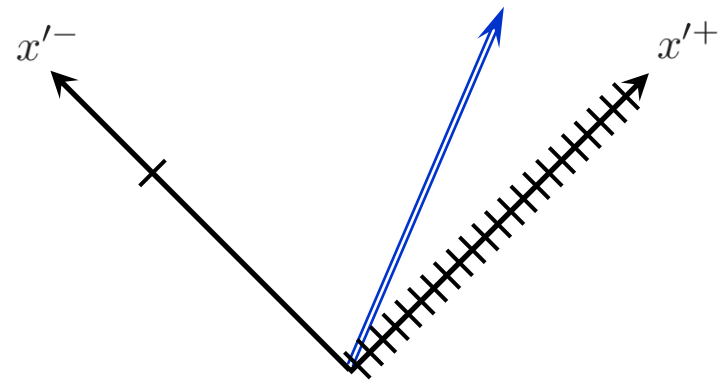
Relativistic (longitudinal) Doppler effect !

Instant form vs light-front form

Infinite-momentum frame $\beta \rightarrow 1$



$$x'^0 \approx x'^3 \approx \frac{x'^+}{\sqrt{2}}$$



$$x'^- \approx 0$$

Galilean/non-relativistic algebra

Galilean symmetry

$$[J^i, J^j] = i\epsilon^{ijk} J^k$$

$$[J^i, B^j] = i\epsilon^{ijk} B^k$$

$$[B^i, B^j] = 0$$

$$[B^i, P^j] = -i\delta^{ij} M$$

$$[B^i, H] = -iP^i$$

$$[B^i, M] = [J^i, M] = [J^i, H] = 0$$

Position operator (center of inertia)

$$B^i = -MR^i$$



$$[R^i, P^j] = i\delta^{ij} \mathbb{1}$$

$$[R^i, R^j] = 0$$

$$[J^i, R^j] = i\epsilon^{ijk} R^k$$

Galilean boost

$$M' = M$$

$$\vec{p}' = \vec{p} + M\vec{v}$$

[Susskind, PR165 (1968) 1535]
[Kogut, Soper, PRD1 (1970) 2901]
[Burkardt, IJMPA18 (2003) 2, 173]
[C.L., EPJC78 (2018) 9, 785]

Light-front Poincaré algebra

$$\begin{aligned} B_{\perp}^1 &= \frac{1}{\sqrt{2}}(K^1 + J^2) & \mathcal{J}_{\perp}^1 &= \frac{1}{\sqrt{2}}(J^1 + K^2) \\ B_{\perp}^2 &= \frac{1}{\sqrt{2}}(K^2 - J^1) & \mathcal{J}_{\perp}^2 &= \frac{1}{\sqrt{2}}(J^2 - K^1) \end{aligned}$$

Transverse space-time symmetry

$$\begin{aligned} [J^3, \mathcal{J}_{\perp}^i] &= i\epsilon^{3ij} \mathcal{J}_{\perp}^j \\ [J^3, B_{\perp}^i] &= i\epsilon^{3ij} B_{\perp}^j \\ [B_{\perp}^i, B_{\perp}^j] &= 0 \end{aligned}$$

$$\begin{aligned} [B_{\perp}^i, P_{\perp}^j] &= -i\delta_{\perp}^{ij} P^+ \\ [B_{\perp}^i, P^-] &= -iP_{\perp}^i \\ [B_{\perp}^i, P^+] &= [J^3, P^+] = [J^3, P^-] = 0 \end{aligned}$$

Transverse position operator (center of P^+)

$$B_{\perp}^i = -P^+ R_{\perp}^i$$



$$\begin{aligned} [R_{\perp}^i, P_{\perp}^j] &= i\delta_{\perp}^{ij} \mathbb{1} \\ [R_{\perp}^i, R_{\perp}^j] &= 0 \\ [J^3, R_{\perp}^i] &= i\epsilon^{3ij} R_{\perp}^j \end{aligned}$$

Transverse boost

$$p'^+ = p^+ \quad \vec{p}'_{\perp} = \vec{p}_{\perp} + p^+ \vec{v}_{\perp}$$

[Susskind, PR165 (1968) 1535]
[Kogut, Soper, PRD1 (1970) 2901]
[Burkardt, IJMPA18 (2003) 2, 173]
[C.L., EPJC78 (2018) 9, 785]

Light-front densities

Localized states

momentum space $|p^+, \vec{p}_\perp\rangle$

mixed space $|p^+, \vec{r}_\perp\rangle = \int \frac{d^2 p_\perp}{(2\pi)^2} e^{-i\vec{p}_\perp \cdot \vec{r}_\perp} |p^+, \vec{p}_\perp\rangle$

Normalizations

$$\begin{aligned} \langle p'^+, \vec{p}'_\perp | p^+, \vec{p}_\perp \rangle &= 2p^+ (2\pi)^3 \delta(p'^+ - p^+) \delta^{(2)}(\vec{p}'_\perp - \vec{p}_\perp) \\ \langle p'^+, \vec{r}'_\perp | p^+, \vec{r}_\perp \rangle &= 2p^+ 2\pi \delta(p'^+ - p^+) \delta^{(2)}(\vec{r}'_\perp - \vec{r}_\perp) \end{aligned}$$

Impact-parameter dependent distributions (IPDs) $x^+ = 0$

$$\begin{aligned} \langle p'^+, \vec{r}'_\perp | \int dx^- J^+(x) | p^+, \vec{r}_\perp \rangle &= 2p^+ 2\pi \delta(p'^+ - p^+) \delta^{(2)}(\vec{r}'_\perp - \vec{r}_\perp) \\ &\times \underbrace{\int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i\vec{\Delta}_\perp \cdot (\vec{x}_\perp - \vec{r}_\perp)} \frac{\langle P^+, \frac{\vec{\Delta}_\perp}{2} | J^+(0) | P^+, -\frac{\vec{\Delta}_\perp}{2} \rangle}{2P^+}}_{\text{Internal distribution}} \end{aligned}$$

NB: $\int dx^- \Leftrightarrow \Delta^+ = 0$ is essential to ensure **probabilistic** interpretation !

Drell-Yan frame

[Soper, PRD15 (1977) 1141]
 [Burkardt, PRD62 (2000) 071503]
 [Diehl, EPJC25 (2002) 223]
 [Burkardt, IJMPA18 (2003) 2, 173]

Light-front densities

Unpolarized quark IPD

$$\vec{b}_\perp = \vec{x}_\perp - \vec{r}_\perp$$

$$q(x, \vec{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i\vec{\Delta}_\perp \cdot \vec{b}_\perp} H_q(x, 0, -\vec{\Delta}_\perp^2)$$

$$H_q(x, \xi, t)$$

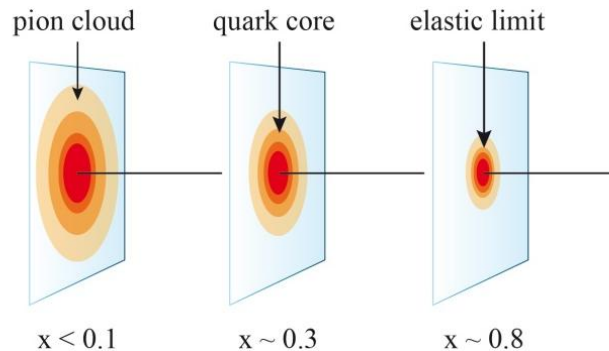
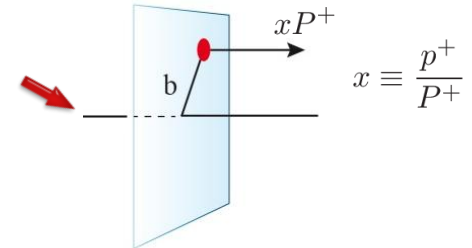
$$\xi = -\frac{\Delta^+}{2P^+}$$

$$t = \Delta^2 = 2\Delta^+ \Delta^- - \vec{\Delta}_\perp^2$$

Center of P^+ (light-front inertia)

$$B_\perp^i = M^{+i} = \int dx^- d^2 x_\perp (x^+ T^{+i} - x_\perp^i T^{++})$$

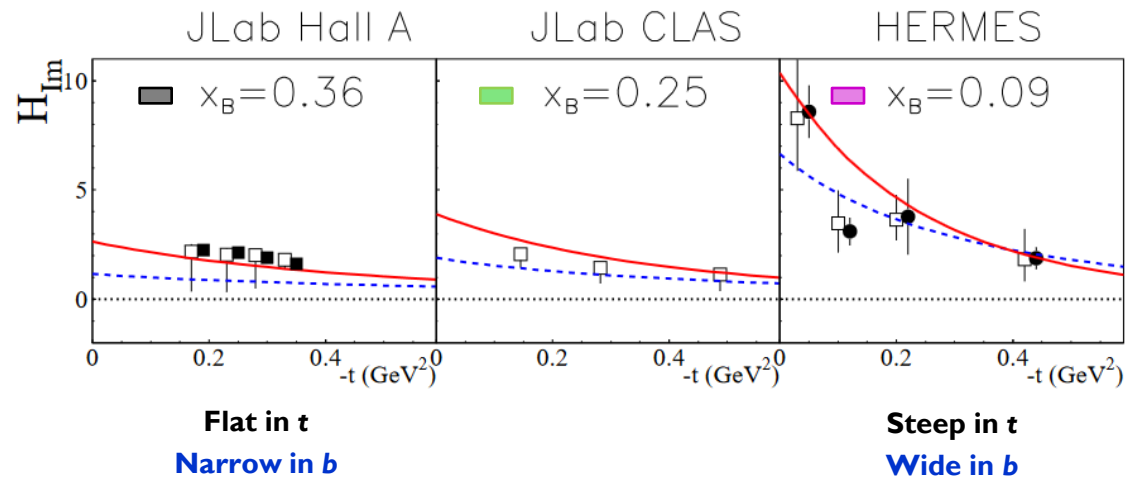
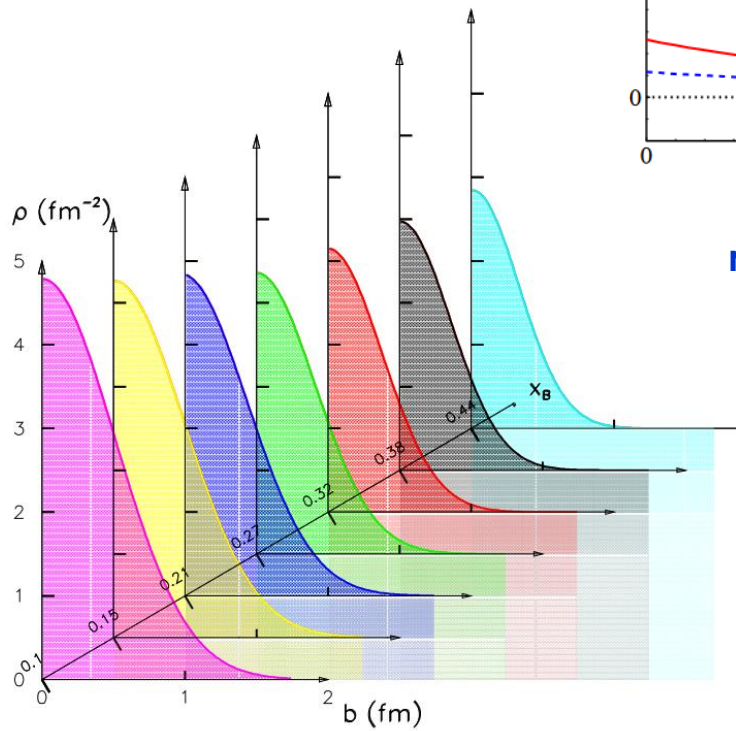
$$\begin{aligned} \vec{R}_\perp &= \frac{1}{P^+} \int dx^- d^2 x_\perp \vec{x}_\perp T^{++}(x) \\ &= \frac{1}{P^+} \sum_n \vec{r}_{\perp,n} p_n^+ = \sum_n x_n \vec{r}_{\perp,n} \end{aligned}$$



$$\langle \vec{b}_\perp^2 \rangle(x) = \frac{\int d^2 b_\perp \vec{b}_\perp^2 q(x, \vec{b}_\perp)}{\int d^2 b_\perp q(x, \vec{b}_\perp)} \xrightarrow{x \rightarrow 1} 0$$

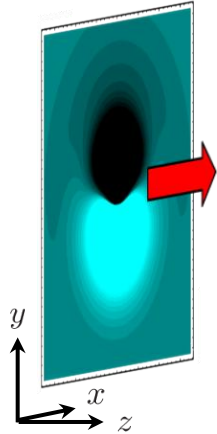
as a result of $x_n \geq 0, \quad \sum_n x_n = 1$

Light-front densities



Light-front artifacts

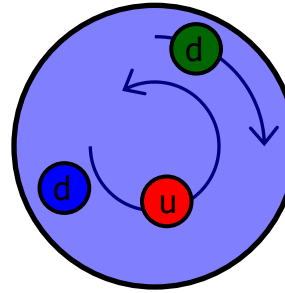
$$J^+ = \frac{J^0 + J^3}{\sqrt{2}}$$



**Neutron
at rest**

$$J_{\text{rest}}^3 < 0$$

$$J_{\text{rest}}^3 > 0$$



$$\vec{S} \otimes$$

$$\mu_n = -1.91$$

**Magnetic dipole
moment**

$$\vec{E}'_{\perp} = \gamma (\vec{E}_{\perp} + \vec{v} \times \vec{B}_{\perp}) \quad \Rightarrow \quad \vec{d}'_{\perp}(\vec{r}') = \gamma \vec{v} \times \vec{\mu}_{\perp}(\vec{r})$$

$$\vec{d}'_{\perp} = \vec{v} \times \vec{\mu}_{\perp}$$

$$\int d^3r' = \int \frac{d^3r}{\gamma}$$

**Induced
electric dipole
moment**

Light-front artifacts

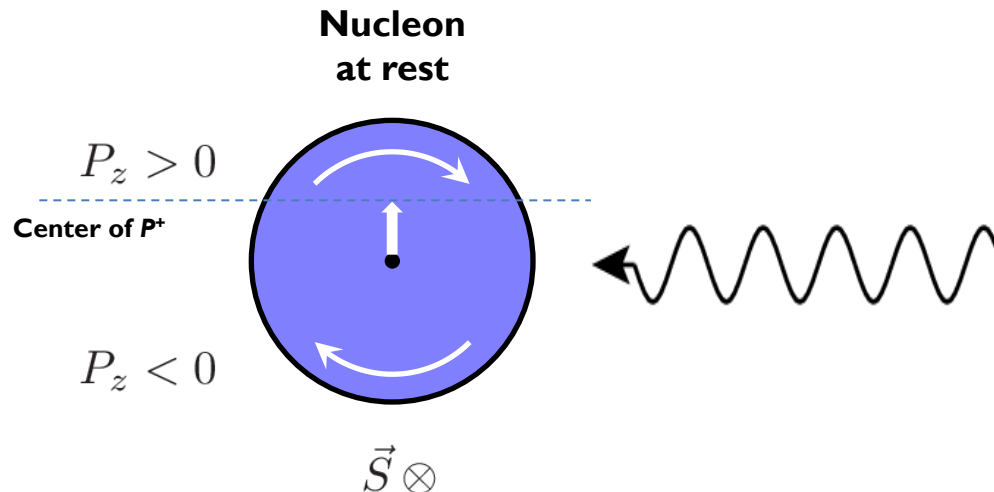
Transverse electric dipole moment

$$\hat{S} = \vec{S}/|\vec{S}|$$

$$\vec{e}_z \Leftrightarrow \vec{v}_{\text{IMF}}$$

$$\langle \vec{d}_\perp \rangle = \int d^2b_\perp \vec{b}_\perp \rho_{\text{LF}}(\vec{b}_\perp) = \frac{\vec{e}_z \times \hat{S}_\perp}{2M_N} \kappa_N$$

$$= \underbrace{\frac{\vec{e}_z \times \hat{S}_\perp}{2M_N} G_M(0)}_{\text{Expected from Lorentz transformation}} - \underbrace{\frac{\vec{e}_z \times \hat{S}_\perp}{2M_N} G_E(0)}_{\substack{\text{Transverse shift} \\ \text{of center of } P^+ \text{ (Electric charge)}}}$$



[Burkardt, IJMPA18 (2003) 2, 173]

[C.L., EPJC78 (2018) 785]

[Chen, C.L., PRD107 (2023) 9, 096003]

Phase-space picture

Position operator in QFT

Textbook claim: « R^μ does not exist »

In reality... it depends on what you mean or want !

Center	Position operator	Canonical relation $[R^i, P^j] = i\delta^{ij}$	Vector under rotations $[J^i, R^j] = i\epsilon^{ijk} R^k$	Compatibility of components $[R^i, R^j] = 0$	Four-vector transformation $R'^\mu = \Lambda^\mu_\nu R^\nu$	
P^+	$R_\perp^i = \frac{1}{P^+} \int dx^- d^2x_\perp x_\perp^i T^{++} = -\frac{B_\perp^i}{P^+}$ $x^+ = 0$	✓	✓	✓	✗	2D
Energy	$R_E^i = \frac{1}{P^0} \int d^3x x^i T^{00} = -\frac{K^i}{P^0}$ $x^0 = 0$	✓	✓	✗	✗	
Mass	$R_M^\mu = \Lambda^\mu_\nu R_E^\nu _{\text{rest}}$	✓	✓	✗	✓	3D
Canonical (or spin)	$R_c^i = \frac{P^0 R_E^i + M R_M^i}{P^0 + M}$	✓	✓	✓	✗	

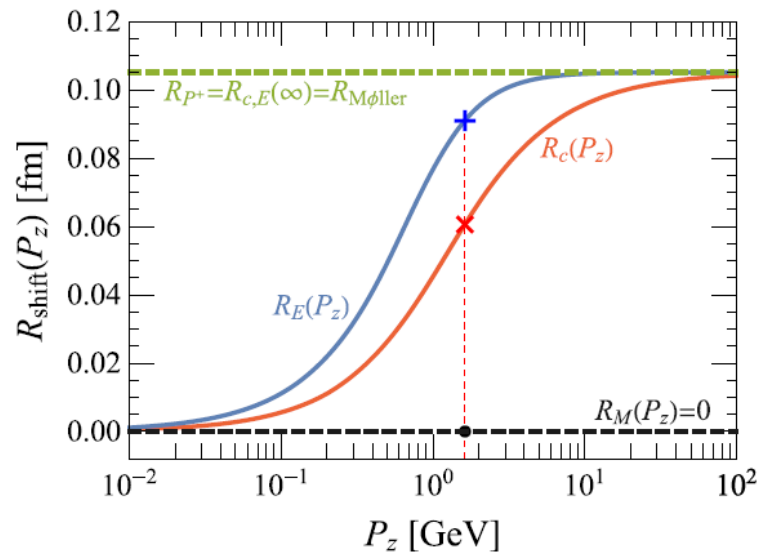
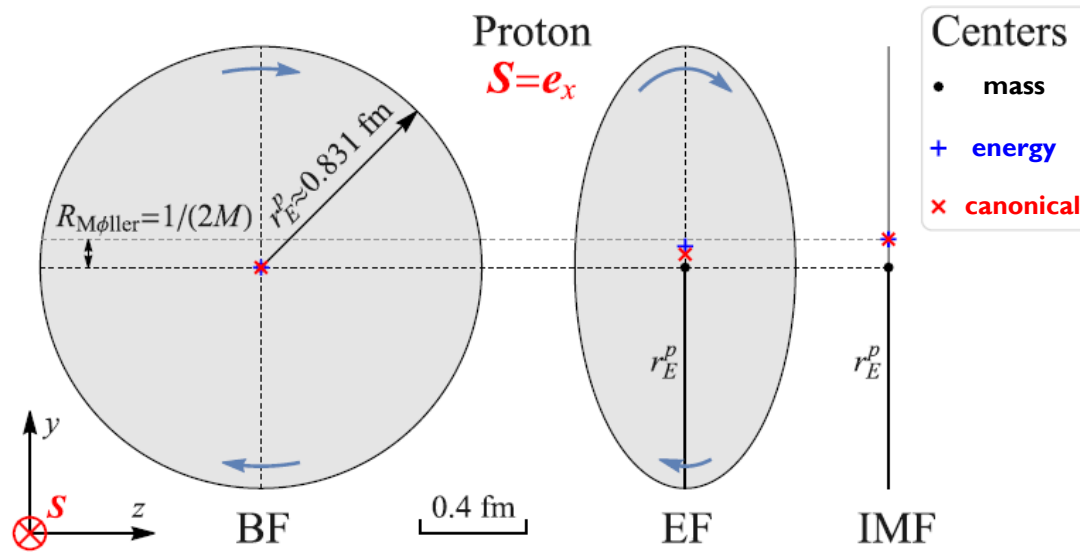
→ Localized states
on spacelike
hypersurface
 $x^0 = 0$

$$\vec{R}_c |\vec{r}\rangle = \vec{r} |\vec{r}\rangle$$

$$|\vec{r}\rangle = \int \frac{d^3p}{(2\pi)^3 \sqrt{2p^0}} e^{-i\vec{p}\cdot\vec{r}} |p\rangle$$

[Pryce, PRSLA195 (1948) 62]
[Newton, Wigner, RMP21 (1949) 3, 400]
[Fleming, PR137 (1965) B188]
[C.L., EPJC78 (2018) 785]

Position operator in QFT



Spatial distributions (general formalism)

Phase-space representation

$$\langle \psi | O(x) | \psi \rangle = \int \frac{d^3 P}{(2\pi)^3} d^3 R \rho_\psi(\vec{R}, \vec{P}) \langle O \rangle_{\vec{R}, \vec{P}}(x)$$

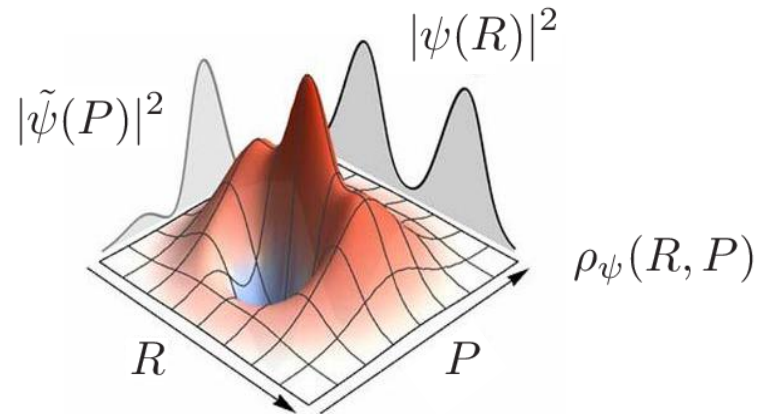
Nucleon Wigner distribution

$$\begin{aligned} \rho_\psi(\vec{R}, \vec{P}) &= \int d^3 z e^{-i\vec{P} \cdot \vec{z}} \psi^*(\vec{R} - \frac{\vec{z}}{2}) \psi(\vec{R} + \frac{\vec{z}}{2}) \\ &= \int \frac{d^3 q}{(2\pi)^3} e^{-i\vec{q} \cdot \vec{R}} \tilde{\psi}^*(\vec{P} + \frac{\vec{q}}{2}) \tilde{\psi}(\vec{P} - \frac{\vec{q}}{2}) \end{aligned}$$

$$\begin{aligned} \psi(\vec{r}) &= \int \frac{d^3 p}{(2\pi)^3} e^{i\vec{p} \cdot \vec{r}} \tilde{\psi}(\vec{p}) \\ \tilde{\psi}(\vec{p}) &= \frac{\langle p | \psi \rangle}{\sqrt{2p^0}} \end{aligned}$$

Quasi-probabilistic interpretation

$$\begin{aligned} \int d^3 R \rho_\psi(\vec{R}, \vec{P}) &= |\tilde{\psi}(\vec{P})|^2 \\ \int \frac{d^3 P}{(2\pi)^3} \rho_\psi(\vec{R}, \vec{P}) &= |\psi(\vec{R})|^2 \end{aligned}$$



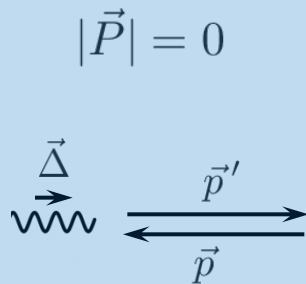
[Wigner, PR40 (1932) 749]
[Carruthers, Zachariasen, PRD13 (1976) 4, 950]
[Hillery, O'Connell, Scully, Wigner, PR106 (1984) 121]
[Bialynicki-Birula, Gornicki, Rafelski, PRD 44 (1991) 1825]

Relativistic spatial distributions

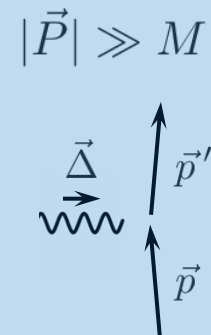
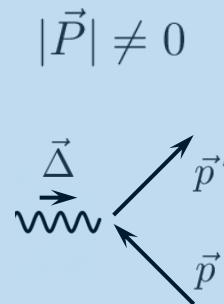
Internal distribution (for a state « localized » in phase-space) $x^0 = 0$

$$\begin{aligned}\langle O \rangle_{\vec{R}, \vec{P}}(\vec{x}) &= \int \frac{d^3\Delta}{(2\pi)^3} e^{-i\vec{\Delta} \cdot (\vec{x} - \vec{R})} \frac{\langle P + \frac{\Delta}{2} | O(0) | P - \frac{\Delta}{2} \rangle}{\sqrt{4(P^0)^2 - (\Delta^0)^2}} \\ &= \langle O \rangle_{\vec{0}, \vec{P}}(\vec{r}), \quad \vec{r} = \vec{x} - \vec{R}\end{aligned}$$

Elastic frames $\Delta^0 = \frac{\vec{P} \cdot \vec{\Delta}}{P^0} \stackrel{!}{=} 0$ (no energy transfer $\Rightarrow$ same initial and final boost factor)



BF



IMF

Relativistic spatial distributions

Fundamental features

- 1) The notion of spatial distribution relies on **simultaneity**
- 2) Probabilistic interpretation requires **factorization** of $\vec{P}$ -dependence

Condition 2) is satisfied by **Galilean symmetry**

... but **not Lorentz symmetry !** $P^0 = \sqrt{\vec{P}^2 + M_N^2(1 + \tau)}$

Traditional perspective: maintain strict probabilistic interpretation by

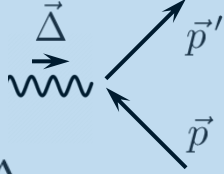
- either neglecting recoil corrections (Sachs approach)
- or restricting to Galilean subgroup (IMF/LF approach)

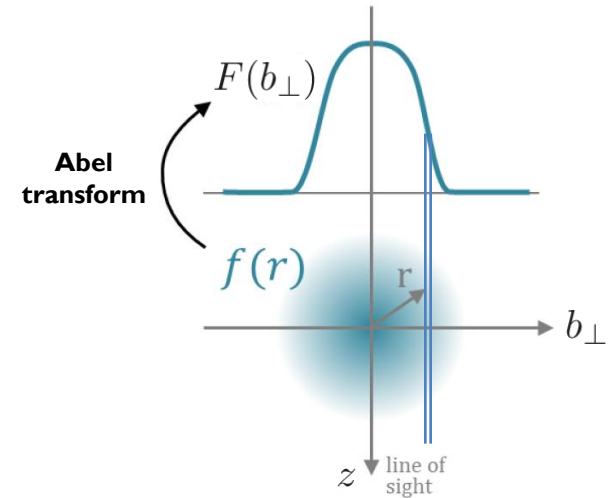
Phase-space perspective: relax to **quasi**-probabilistic interpretation
but fully account for frame dependence !

Relativistic spatial distributions

Elastic frame

$$\vec{P} = P_z \vec{e}_z \Rightarrow \Delta^0 = \frac{P_z \Delta_z}{P^0}$$

$$\Delta^0 = 0 \Rightarrow \Delta_z = 0 \Leftrightarrow \int dz$$




2D charge distribution

$$\rho_E^{\text{EF}}(\vec{b}_\perp; P_z) \equiv \int dz \langle J^0 \rangle_{\vec{R}, P_z \vec{e}_z}(\vec{x})$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i \vec{\Delta}_\perp \cdot \vec{b}_\perp} \frac{\langle p', s' | J^0(0) | p, s \rangle}{2P^0} \Big|_{\text{EF}}$$

Interpolates between BF and IMF

$$\rho_E^{\text{EF}}(\vec{b}_\perp; 0) = \int dz \rho_E^{\text{BF}}(\vec{r})$$

$$\rho_E^{\text{EF}}(\vec{b}_\perp; \infty) = \rho_E^{\text{IMF}}(\vec{b}_\perp)$$

$$\vec{b}_\perp = \vec{x}_\perp - \vec{R}_\perp$$

[C.L., Mantovani, Pasquini, PLB776 (2018) 38]

[C.L., PRL125 (2020) 232002]

[Panteleeva, Polyakov, PRD104 (2021) 1, 014008]

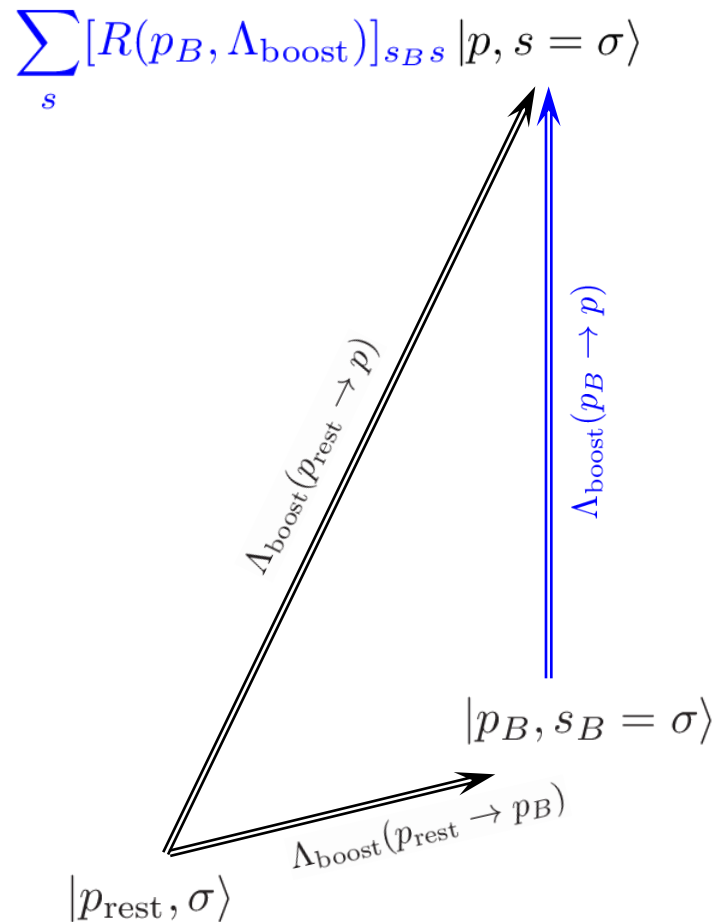
[Kim, Kim, PRD104 (2021) 7, 074003]

Thomas-Wigner rotation



Relativistic boosts do not commute !

$$[K^i, K^j] = -i\epsilon^{ijk} J^k$$



[Thomas, Nature 117 (1926) 514]
[Wigner, ZP124 (1948) 665]
[Wigner, RMP29 (1957) 255]

Four-current amplitude

Lorentz transformation of an off-forward amplitude

$$\langle p', s' | J^\mu(0) | p, s \rangle = \sum_{s'_B, s_B} D_{s'_B s'}^{*(j)}(p'_B, \Lambda) D_{s_B s}^{(j)}(p_B, \Lambda) \Lambda^\mu{}_\nu \langle p'_B, s'_B | J^\nu(0) | p_B, s_B \rangle$$

[Durand, De Celles, Marr, PR126 (1962) 1882]

Confirmed by explicit evaluation of spinors

$$\hat{\Delta} = \vec{\Delta}/|\vec{\Delta}|$$

$$Q^2|_{\text{EF}} = \vec{\Delta}_\perp^2$$

$$\begin{aligned} \langle p', s' | J^0(0) | p, s \rangle|_{\text{EF}} = 2M_N \gamma \Big\{ & \left[\delta_{s's} \cos \theta + (\vec{\sigma}_{s's} \times i\hat{\Delta}_\perp)_z \sin \theta \right] G_E(Q^2) \\ & + \beta \left[-\delta_{s's} \sin \theta + (\vec{\sigma}_{s's} \times i\hat{\Delta}_\perp)_z \cos \theta \right] \sqrt{\tau} G_M(Q^2) \Big\} \end{aligned}$$

$$\begin{aligned} \langle p', s' | J^3(0) | p, s \rangle|_{\text{EF}} = 2M_N \gamma \Big\{ & \beta \left[\delta_{s's} \cos \theta + (\vec{\sigma}_{s's} \times i\hat{\Delta}_\perp)_z \sin \theta \right] G_E(Q^2) \\ & + \left[-\delta_{s's} \sin \theta + (\vec{\sigma}_{s's} \times i\hat{\Delta}_\perp)_z \cos \theta \right] \sqrt{\tau} G_M(Q^2) \Big\} \end{aligned}$$

$$\langle p', s' | J_\perp^i(0) | p, s \rangle|_{\text{EF}} = 2M_N (\sigma_z)_{s's} (\vec{e}_z \times i\hat{\Delta}_\perp)^i \sqrt{\tau} G_M(Q^2)$$

Boost parameters

$$\beta = \frac{P_z}{P^0} \quad \gamma = \frac{P^0}{\sqrt{P^2}}$$

$$P^2 = M_N^2(1 + \tau)$$

Wigner rotation

$$\begin{aligned} \cos \theta &= \frac{P^0 + M_N(1 + \tau)}{(P^0 + M_N)\sqrt{1 + \tau}} \\ \sin \theta &= -\frac{\sqrt{\tau} P_z}{(P^0 + M_N)\sqrt{1 + \tau}} \end{aligned}$$

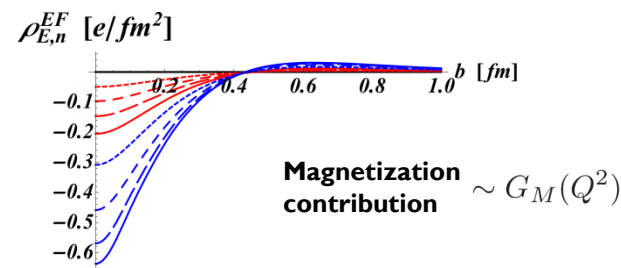
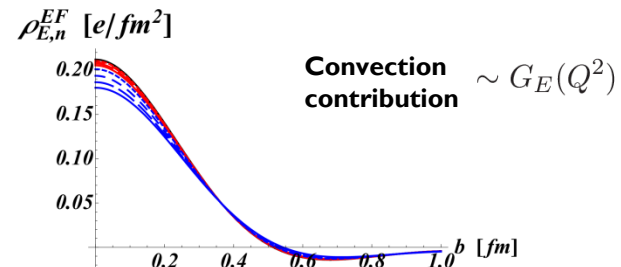
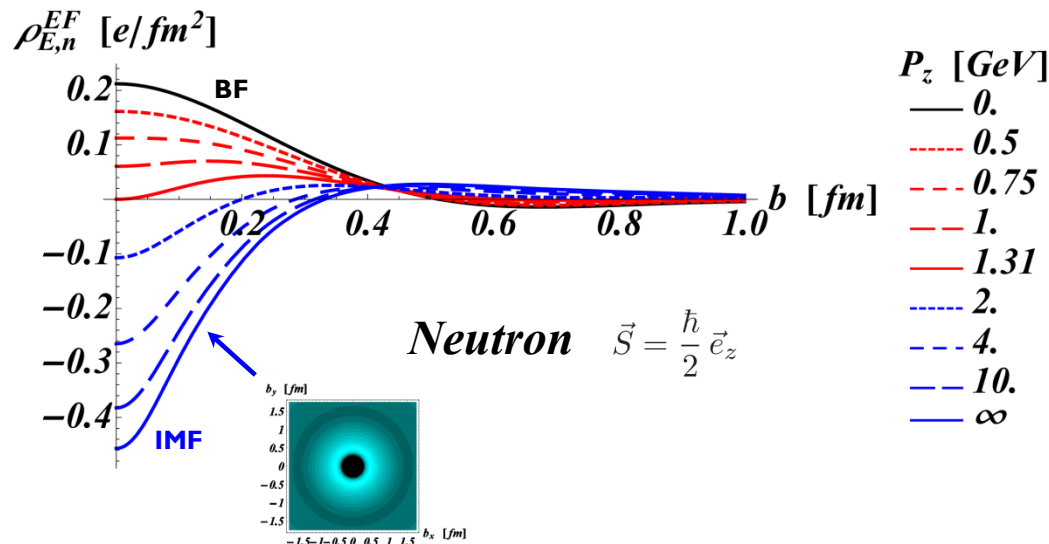
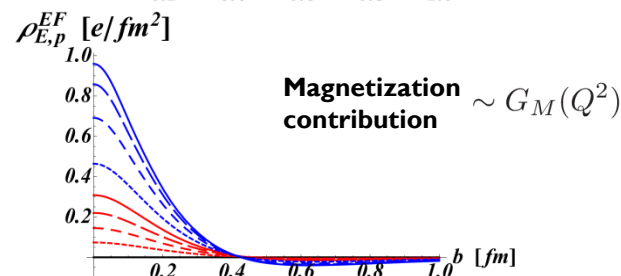
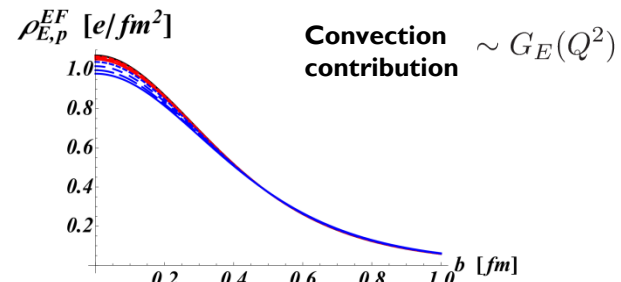
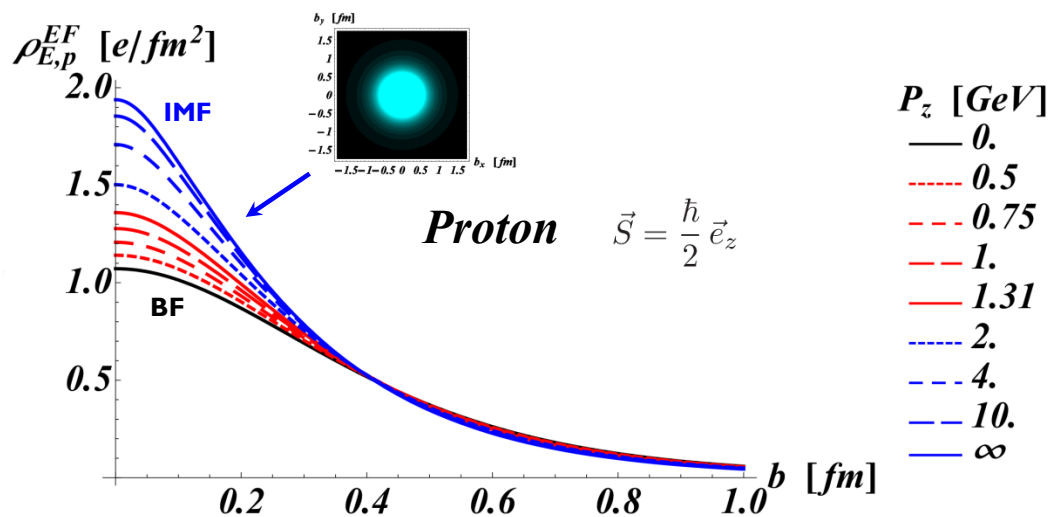
[Chung, Polyzou, Coester, Keister, PRC37 (1988) 2000]

[Rinehimer, Miller, PRC80 (2009) 015201]

[C.L., Wang, PRD105 (2022) 9, 096032]

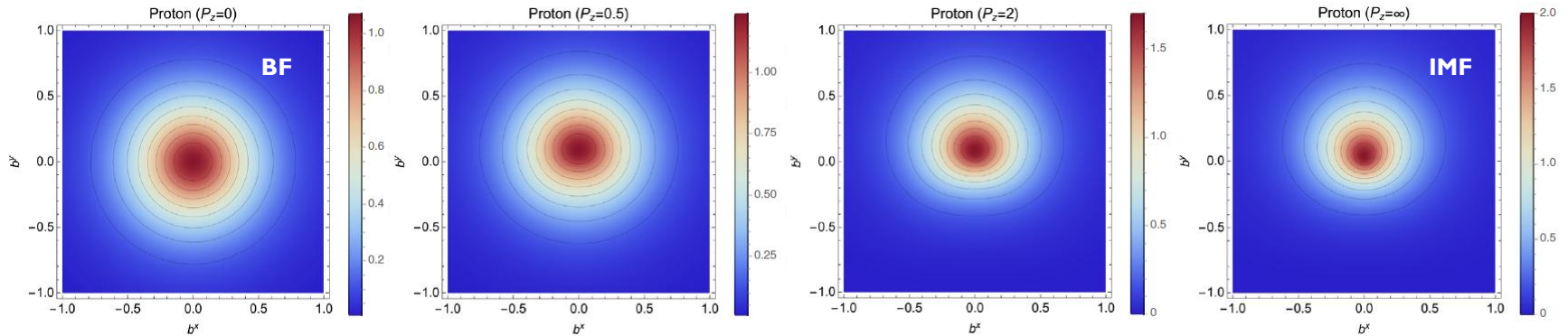
[Chen, C.L., PRD106 (2022) 11, 116024]

EF charge distributions (longitudinal polarization)



EF charge distributions (transverse polarization)

Proton $\vec{S} = \frac{\hbar}{2} \vec{e}_x$



Neutron $\vec{S} = \frac{\hbar}{2} \vec{e}_x$

